



School of Modern Optics

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Lecture 5

Optomechanics of fluid interfaces

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www.loma.cnrs.fr/spip.php?rubrique331

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Introduction

Per photon of a light beam

Energy

$$E = \hbar\omega$$

Linear momentum

$$\mathbf{p} = \hbar\mathbf{k}$$

Angular momentum

$$j_z = s_z + l_z$$

Mechanical principles

Forces

$$\frac{d\mathbf{P}}{dt} = \sum \mathbf{F}_{ext}$$

Torques

$$\frac{d\mathbf{J}}{dt} = \sum \mathbf{\Gamma}_{ext}$$

Mechanical consequence : **light can rotate matter**

Introduction

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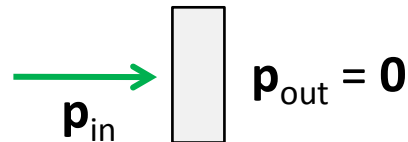
Mechanical consequence : **light can displace matter**

Introduction : optical forces

$$\frac{d\mathbf{P}}{dt} = \sum \mathbf{F}_{ext}$$

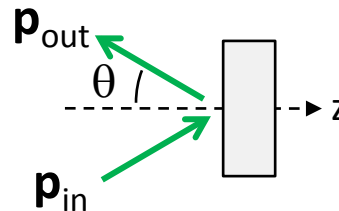
Transferred momentum
per photon

Absorption



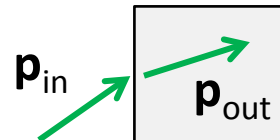
$$\hbar \mathbf{k}_{in}$$

Reflexion



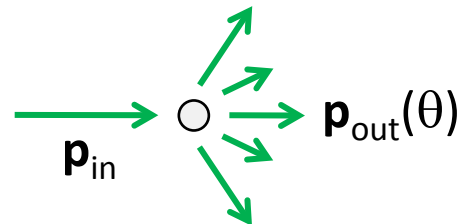
$$2\hbar \cos\theta k_{in} z$$

Refraction



$$\hbar(\mathbf{k}_{in} - \mathbf{k}_{out})$$

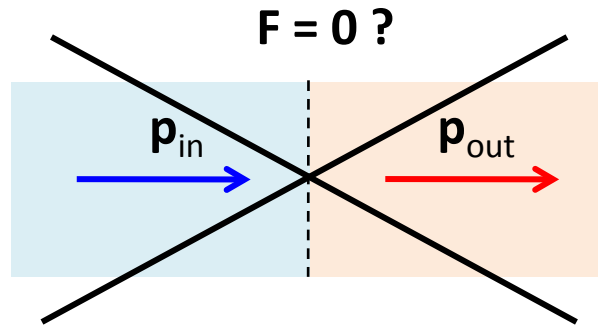
Scattering



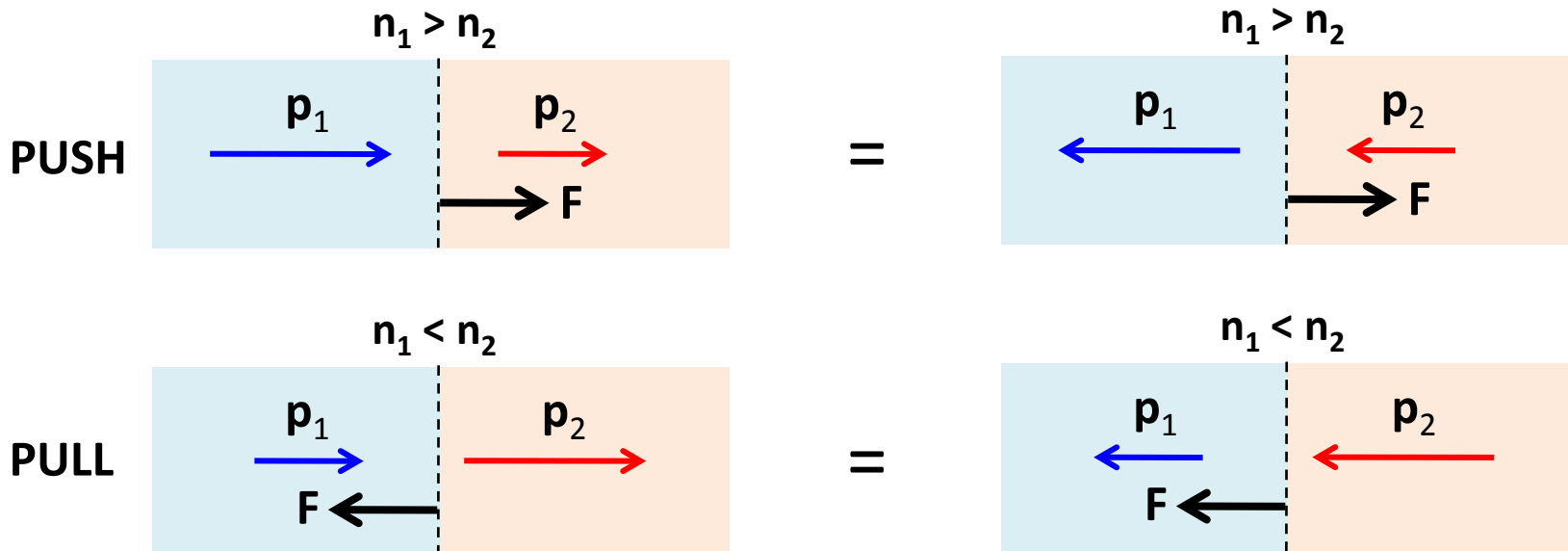
$$\hbar \mathbf{k}_{in} - \hbar \int f(\theta) \mathbf{k}_{out}(\theta) d\theta$$

Redirection

Introduction : optical forces

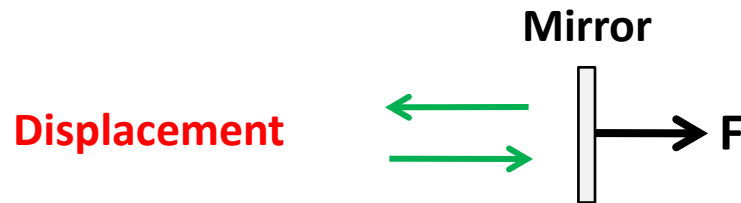


A photon in a dielectric medium : $\mathbf{p} = n\hbar\mathbf{k}$



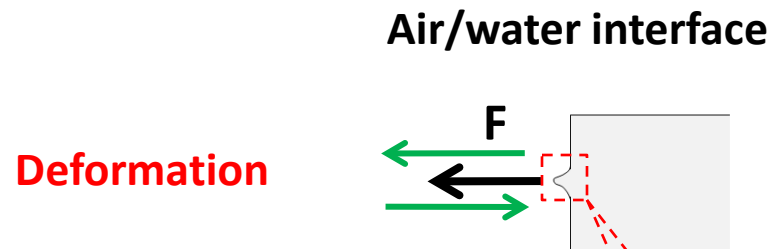
Optical force per unit surface : **radiation pressure of light**

Introduction : optical radiation pressure in the lab



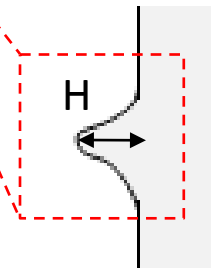
P. Lebedev, Ann. Phys. **6**, 433 (1901)

E.F. Nichols and G.F. Hull, Phys. Rev. **13**, 307 (1901)



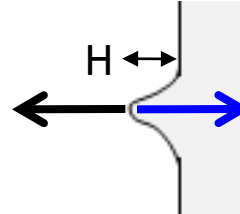
A. Ashkin and J.M. Dziedzic, PRL **30**, 139 (1973)

How large is the effect ?



Introduction : estimation of the fluid elevation

Equilibrium of the interface



$$\Pi_{\text{radiation}} = \Pi_{\text{Laplace}}$$

$$\Pi_{\text{radiation}} = F / S$$

$$= \Phi(n_1 - n_2)\hbar k$$

$$= \Phi(n_1 - n_2) \frac{\hbar\omega}{c}$$

$$= (n_1 - n_2) \frac{I}{c}$$

Number of photons
per unit time and surface

$$\Pi_{\text{Laplace}} = \sigma \kappa$$

$$\approx \frac{\sigma H}{w^2}$$

Interface curvature

Surface tension

Soft interfaces are desirable

$$H \approx (n_1 - n_2) \frac{I w^2}{\sigma c} \approx (n_1 - n_2) \frac{P}{\sigma c}$$

$$P = 1 \text{ W} , \sigma = 10^{-3} - 10^{-2} \text{ N/m} \Rightarrow H \approx 10 - 100 \text{ nm}$$

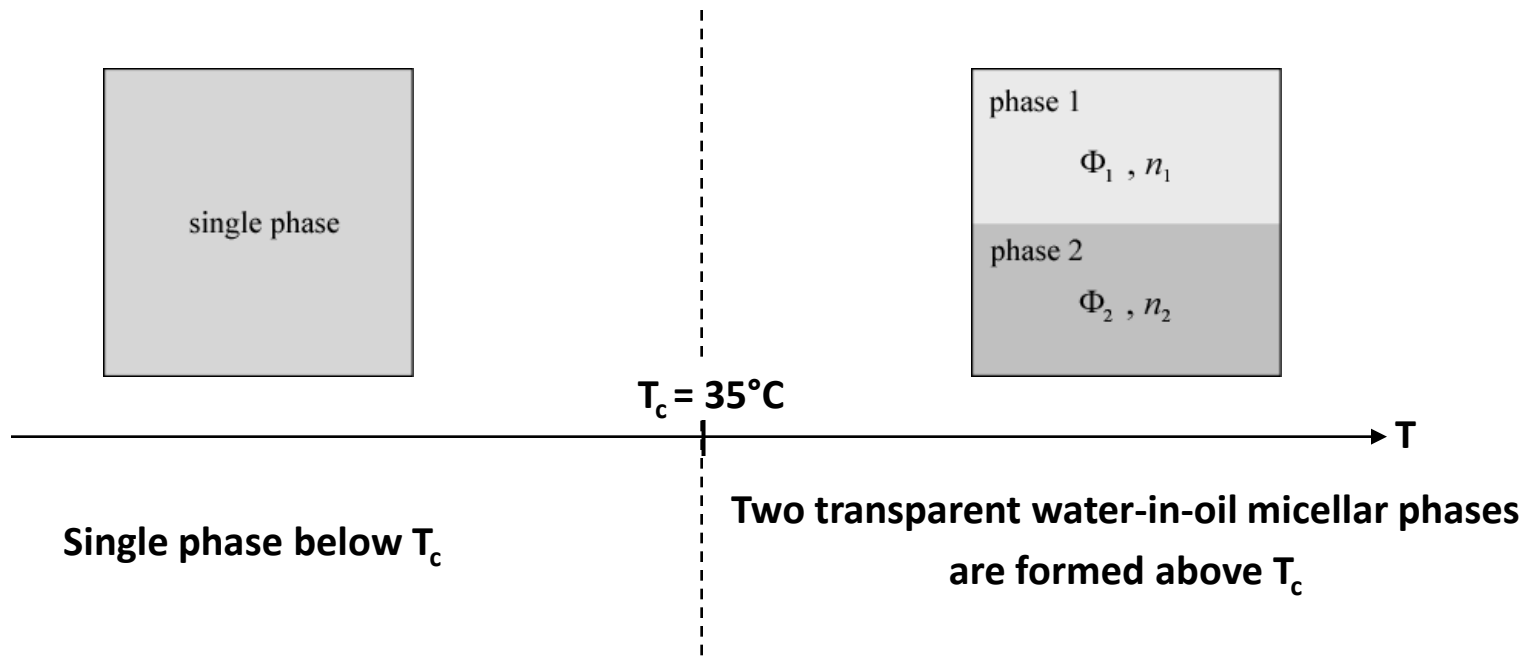
Outline

- 1. Laser induced giant fluid interface deformations**
2. Liquid-core liquid-cladding self-induced optical fibers
3. The optical micro-pipeline
4. Acoustic analogies

A very soft interface in the lab

Quaternary liquid mixture

(toluene, butanol, water, surfactant)

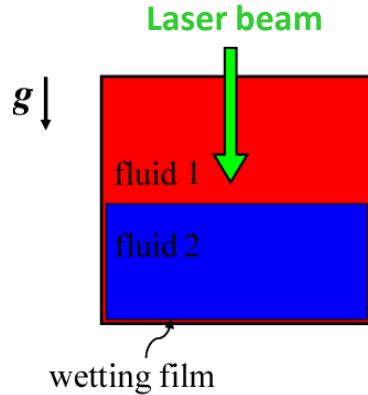


Interfacial surface tension vanishes at $T = T_c$

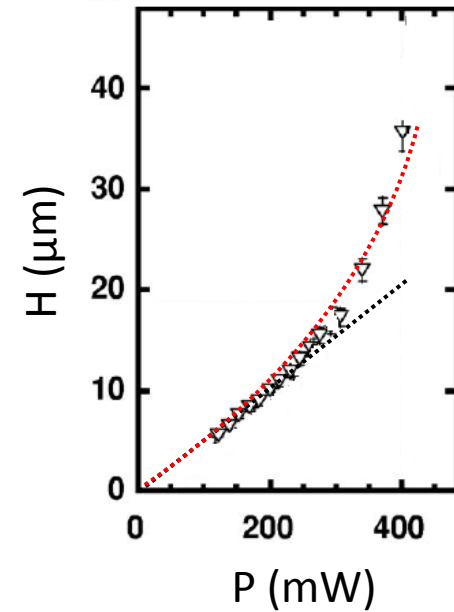
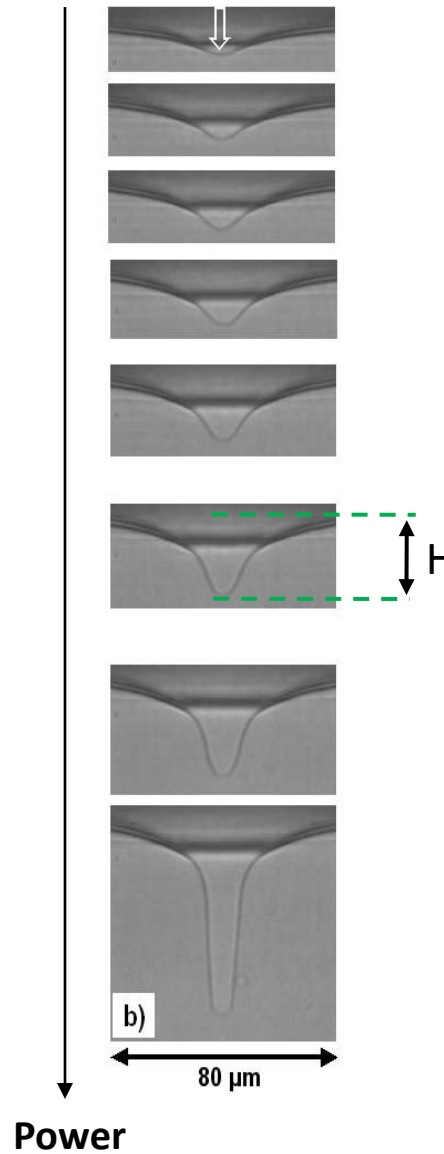
$$\sigma \rightarrow 0$$

($\sigma \sim 10^{-7} - 10^{-6}$ N/m for $\Delta T = T - T_c \sim 2$ to 15K)

Giant fluid interface deformations by light

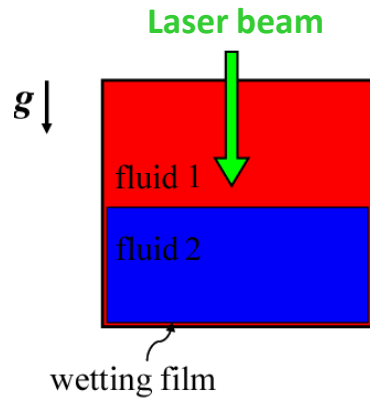


A. Casner and J.P. Delville, PRL **87**, 054503 (2001)

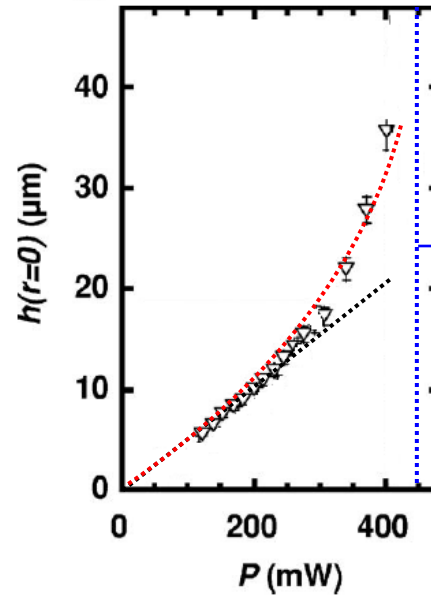


Nonlinear regime is observed

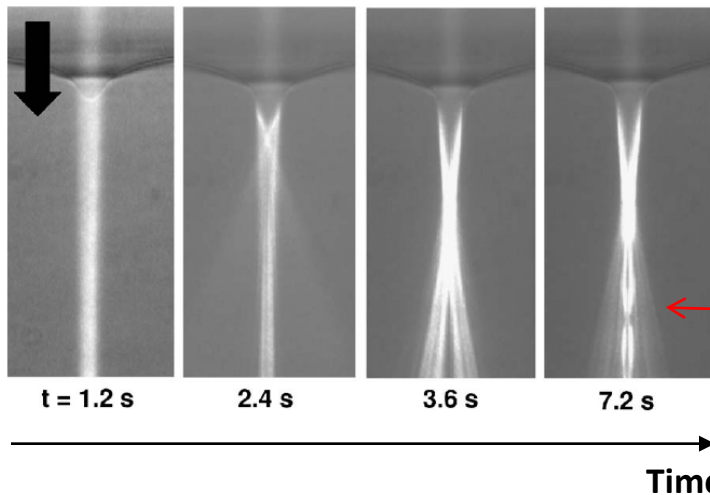
Fluid interface instability induced by light



A. Casner and J.P. Delville, PRL **90**, 144503 (2003)



Jetting transient dynamics



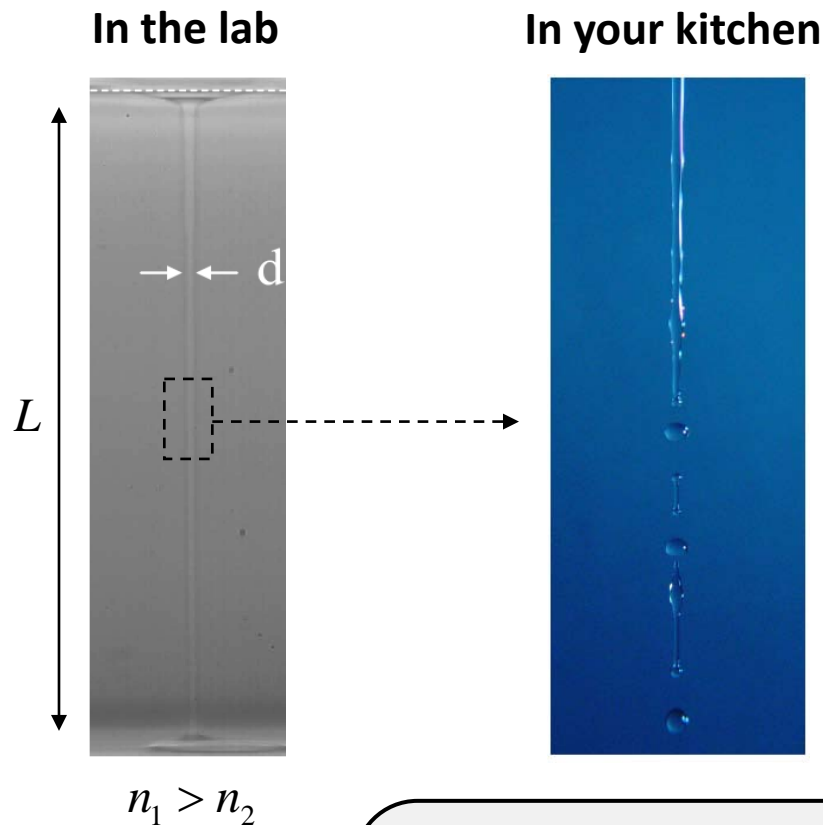
Step-index optical fiber

Total internal reflection

Outline

1. Laser induced giant fluid interface deformations
- 2. Liquid-core liquid-cladding self-induced optical fibers**
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Existence of a liquid fiber : a surprising effect



Why the liquid cylinder
does not collapse (as usual) ?



Capillary effects balanced by light !

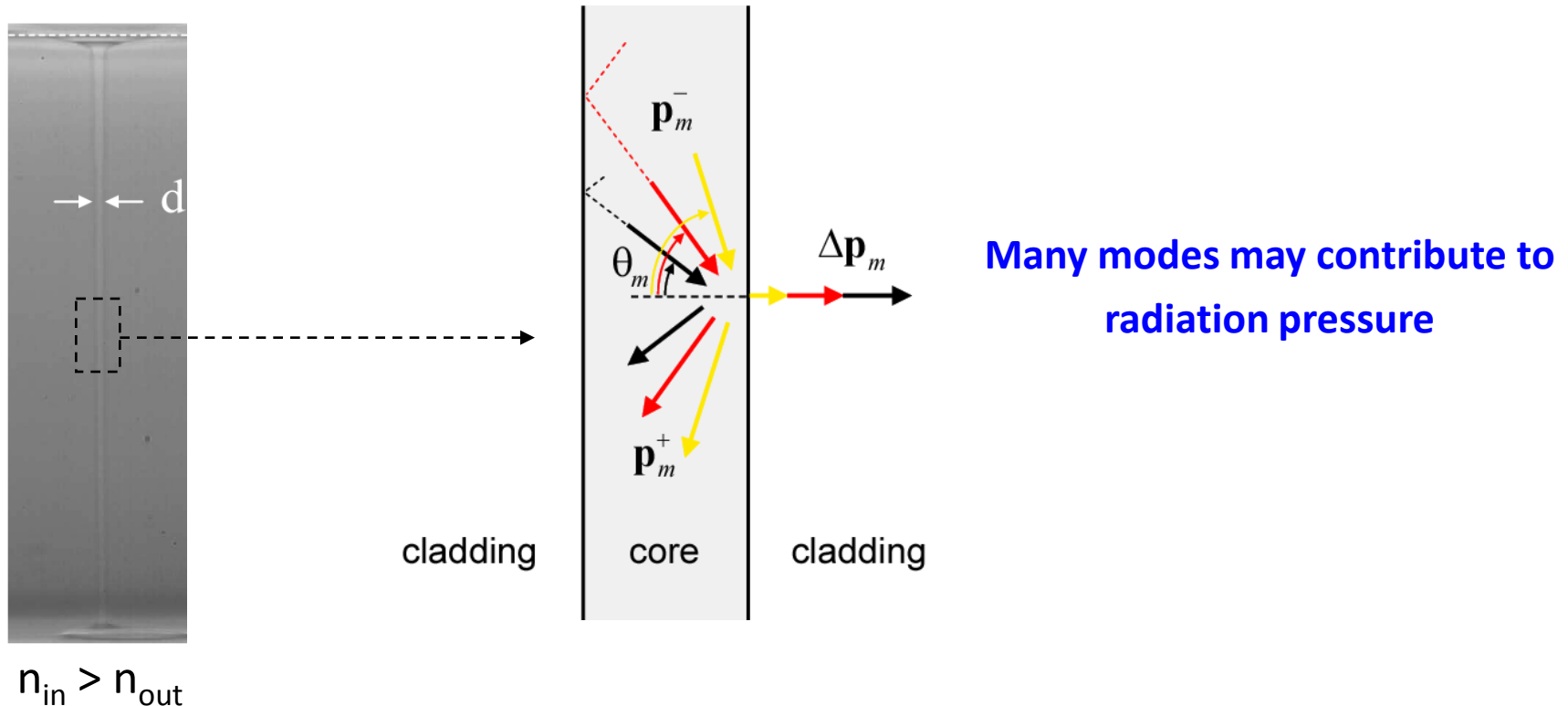
$$\Lambda = L / d > \pi$$

Rayleigh-Plateau instability

A cylindrical liquid column with fixed ends and volume breaks into droplets when the aspect ratio is large enough.

Existence of a liquid fiber : qualitative considerations

Liquid-core liquid-cladding optical fiber sustained by
light radiation pressure



$$\Pi_{\text{radiation}} = \sum_m \Phi_m \Delta \mathbf{p}_m \cdot \mathbf{u}_r$$

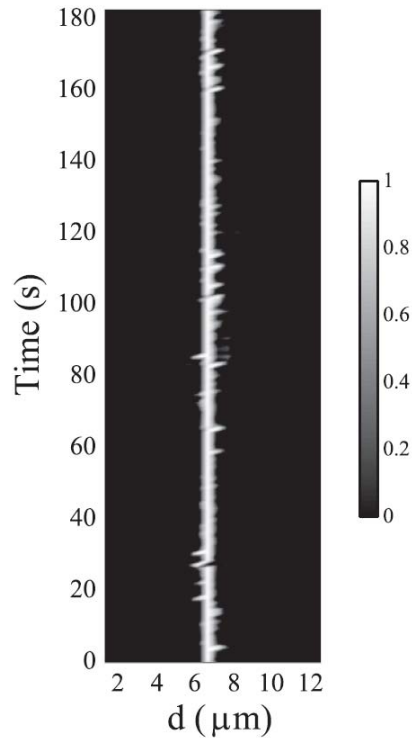
Ray optics approach

Step-index liquid fiber : experiment

Laser beam

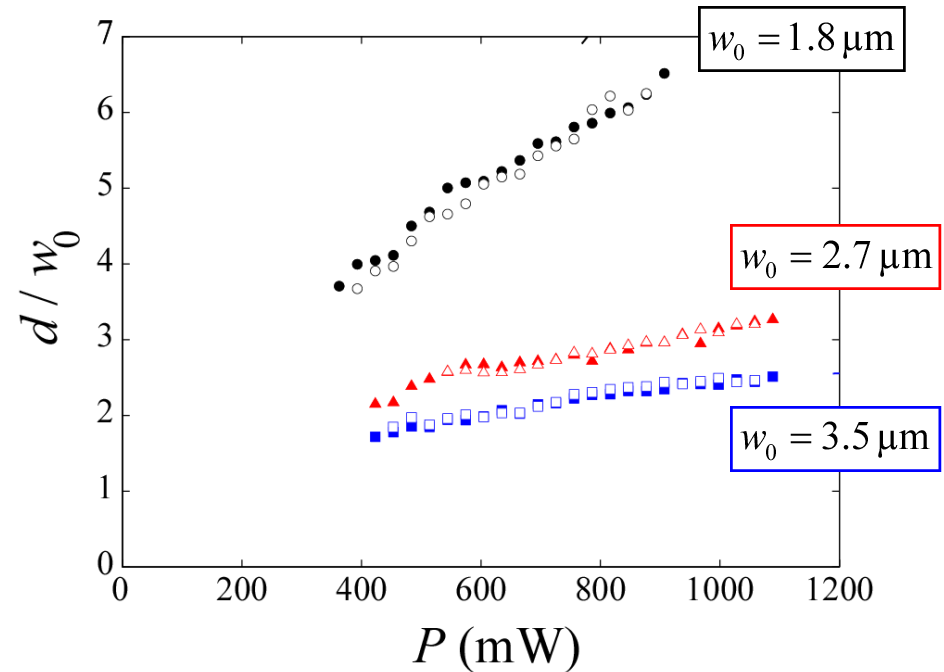


Histogram vs. time



Single-valued and stationary
diameter distribution

Dependence on incident power and waist



Self-adapted liquid optical fiber

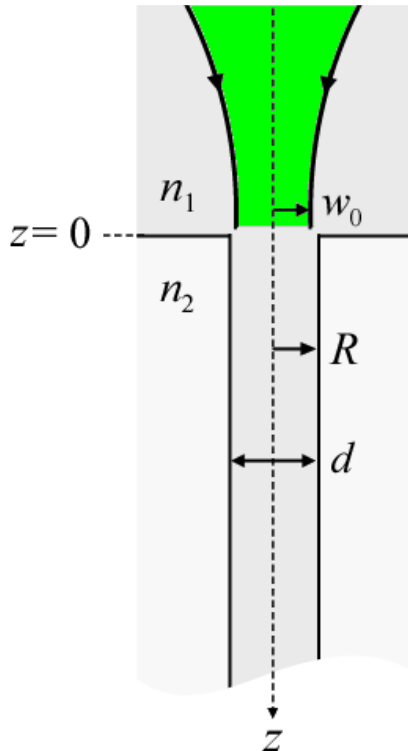
E. Brasselet *et al*, PRL **101**, 014501 (2008)

Step-index liquid fiber : model

Equilibrium equation

$$\Pi_{\text{Laplace}} = \Pi_{\text{radiation}} \quad \begin{cases} \Pi_{\text{Laplace}} = \frac{\sigma}{R} \\ \Pi_{\text{radiation}} = \epsilon_0 \frac{n_1^2 - n_2^2}{4} \left(|\mathbf{E}_{t,1}|^2 + \frac{n_1^2}{n_2^2} |\mathbf{E}_{n,1}|^2 \right)_{r=R} \end{cases}$$

Incident laser beam



Weakly guiding fiber approximation

$$\mathbf{E}^{(m)} = E_0^{(m)} \mathcal{R}_m(r) e^{i\beta_m z} \mathbf{u}_x$$

$$\Downarrow$$

$$\Pi_{\text{radiation}} = \frac{1}{2} \epsilon_0 n (n_1 - n_2) \sum_m \frac{|E_0^{(m)}|^2}{\text{blue arrow pointing to } m}$$

Depends on power transmitted to mode m (LP_{0m})

$$P_m = T_m P$$

$$T_m = \frac{\left| \int_0^\infty \mathbf{E}^{(m)} \cdot \mathbf{E}^{(\text{inc})} r dr \right|^2}{\int_0^\infty |\mathbf{E}^{(m)}|^2 r dr \int_0^\infty |\mathbf{E}^{(\text{inc})}|^2 r dr}$$

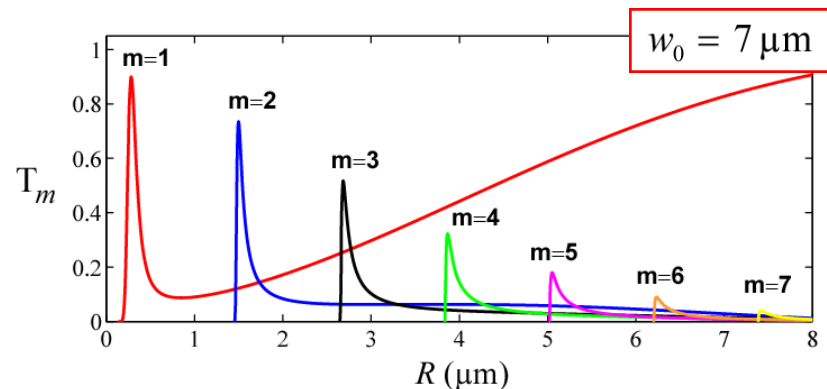
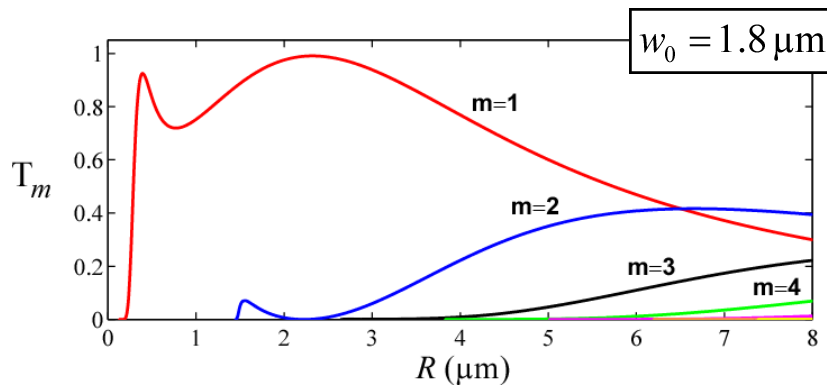
Step-index liquid fiber : model

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Modal transmissions



Weakly guiding fiber approximation

$$\mathbf{E}^{(m)} = E_0^{(m)} \mathcal{R}_m(r) e^{i\beta_m z} \mathbf{u}_x$$



$$\Pi_{\text{radiation}} = \frac{1}{2} \epsilon_0 n (n_1 - n_2) \sum_m \frac{|E_0^{(m)}|^2}{\text{}} \quad \swarrow$$

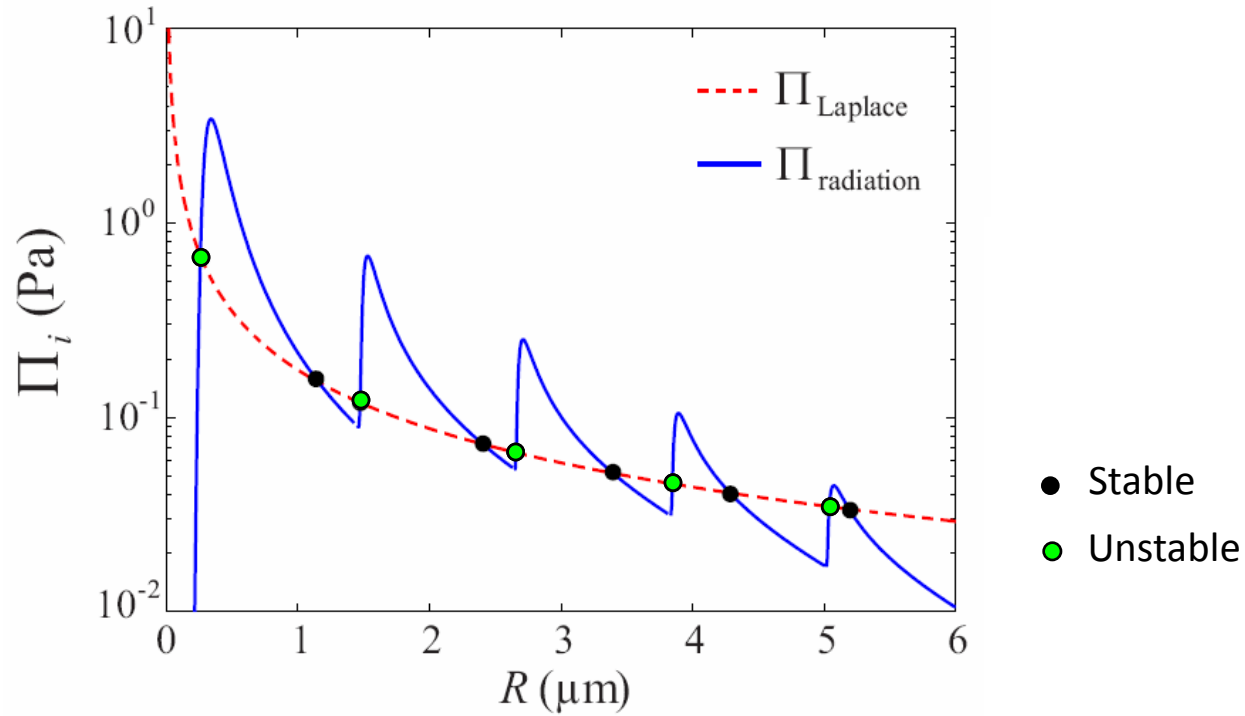
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Step-index liquid fiber : model

$w_0 = 7\mu\text{m}$
 $P = 500\text{mW}$
 $\Delta T = 2\text{K}$



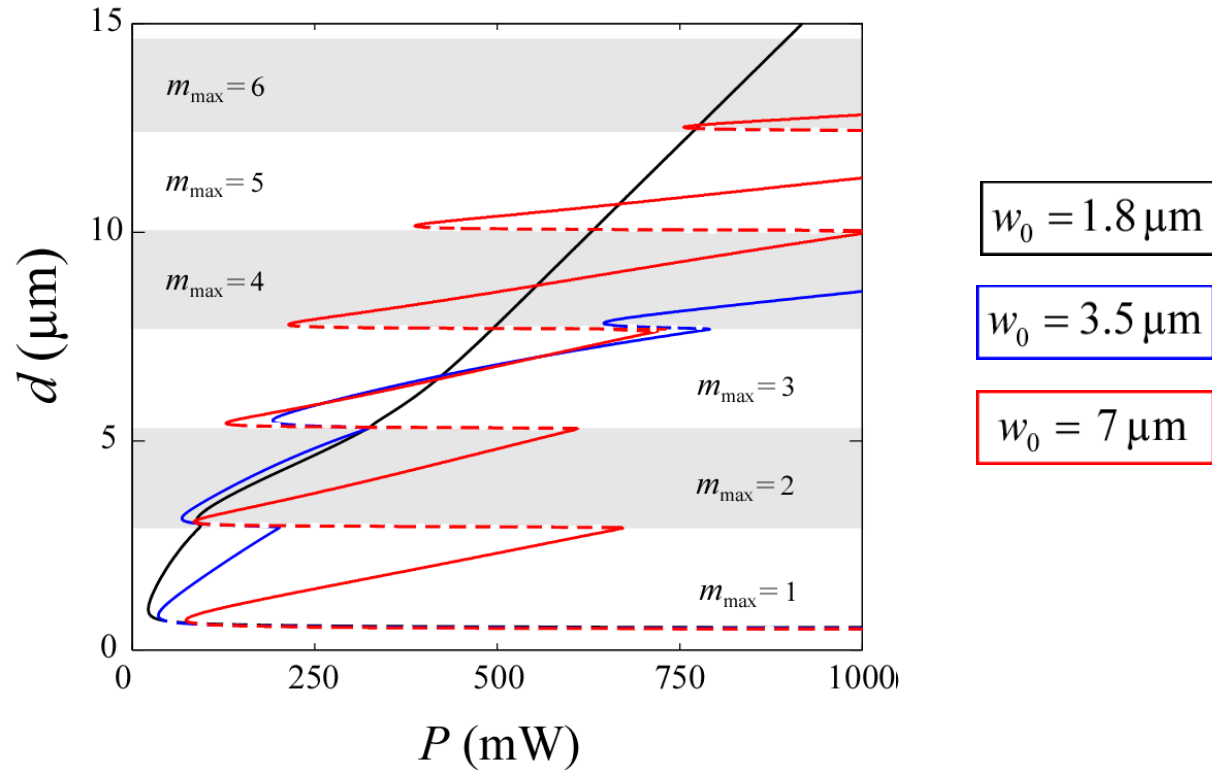
Equilibrium states

$$\Pi_{\text{Laplace}} = \Pi_{\text{radiation}}$$

Stability criterion

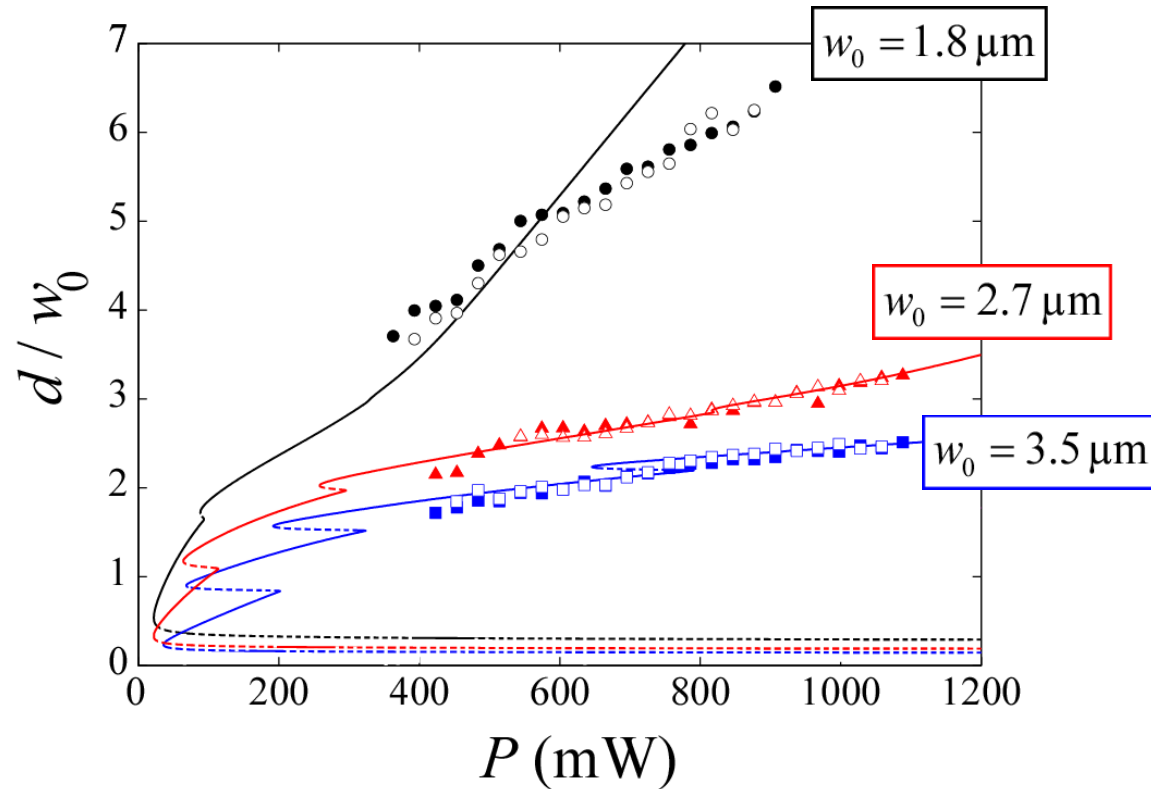
$$\left. \frac{\partial \Pi_{\text{radiation}}}{\partial R} \right|_{R=R_{\text{eq}}^{(n)}} < \left. \frac{\partial \Pi_{\text{Laplace}}}{\partial R} \right|_{R=R_{\text{eq}}^{(n)}}$$

Step-index liquid fiber : model



The core diameter can be either single or multivalued
and sustained by one or several guided modes

Step-index liquid fiber : experiment / model



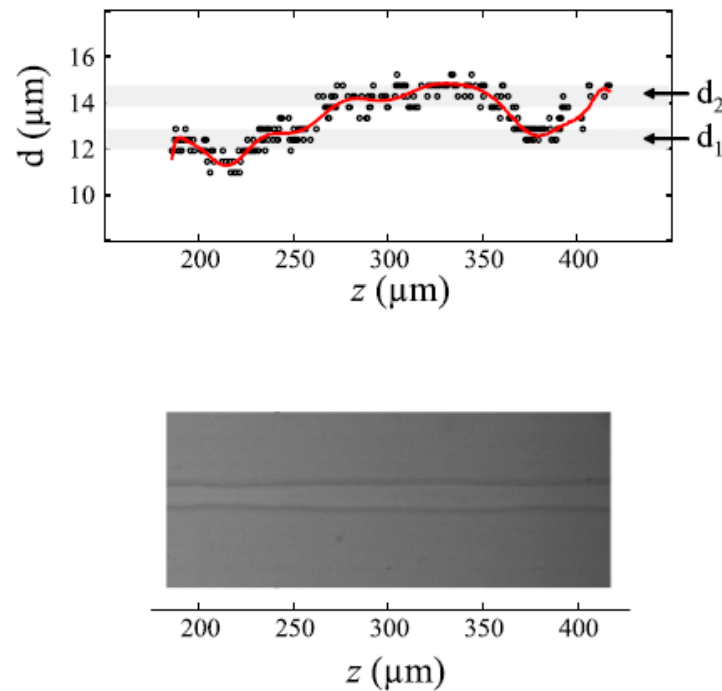
Satisfying agreement at lower beam waists

Step-index liquid fiber : qualitative multistable core behavior

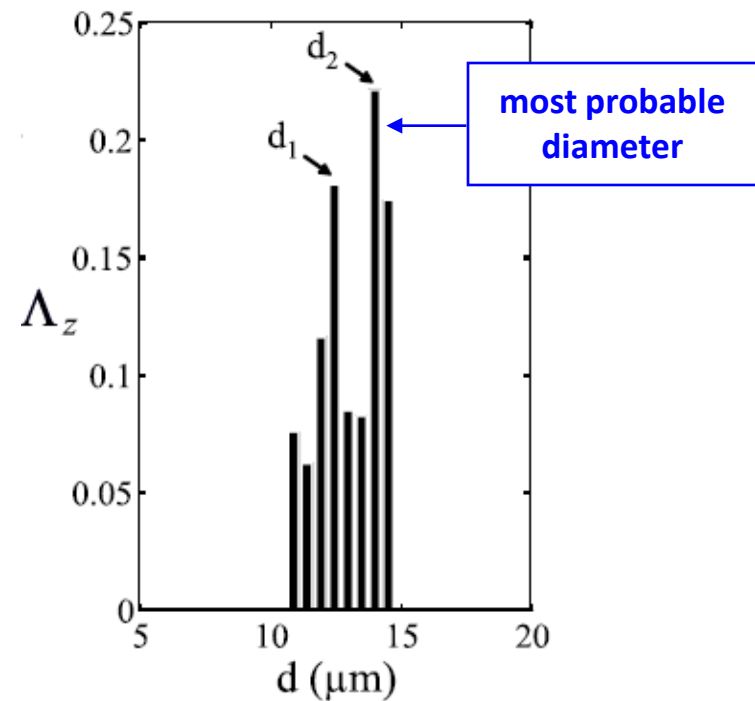
Signature of multistability at larger beam waists **at a fixed time**

$$w_0 = 7.0 \mu\text{m}$$

$$P = 575 \text{ mW}$$



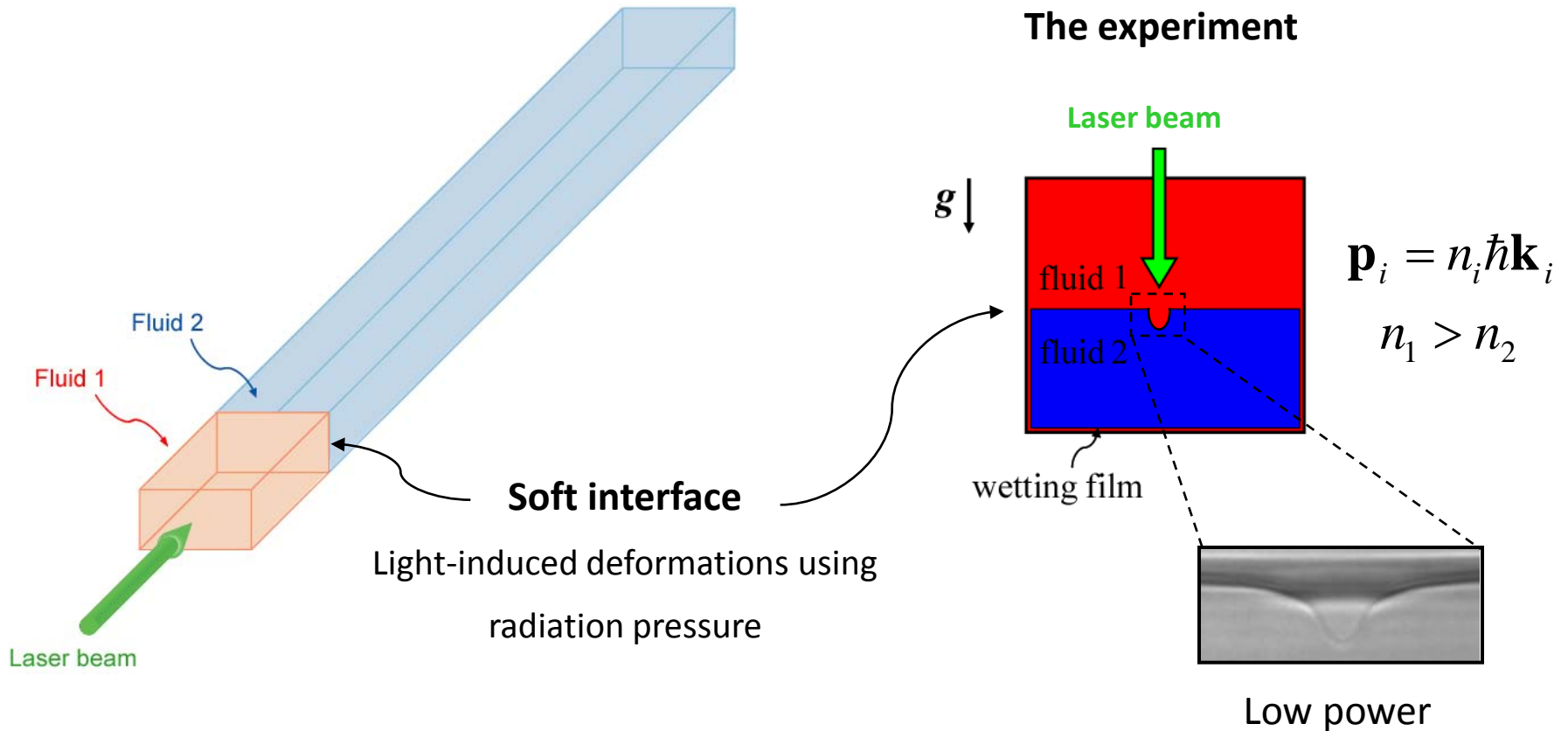
Diameter probability distribution



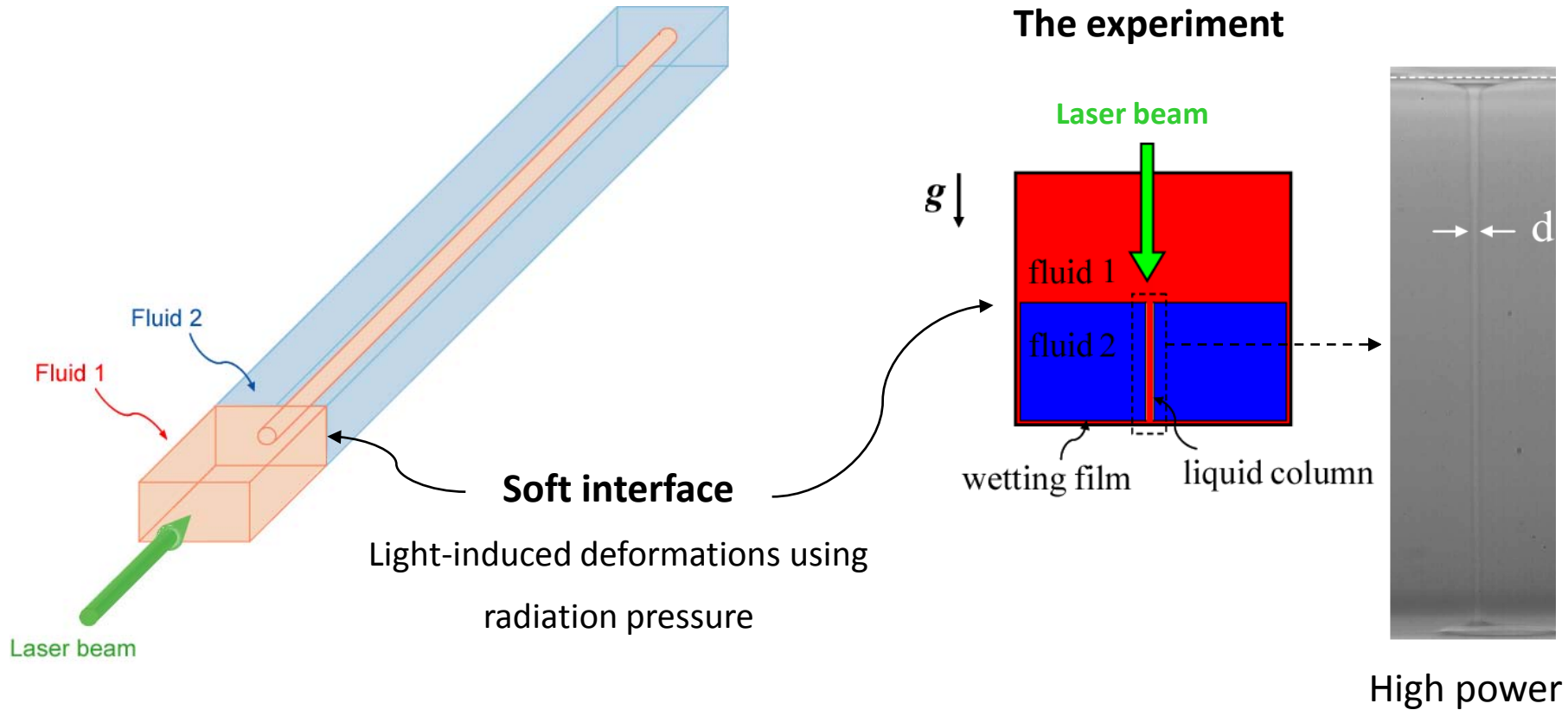
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- 3. The optical micro-pipeline**
4. Acoustic analogies

All-liquid micro-optical pipeline concept



All-liquid micro-optical pipeline concept

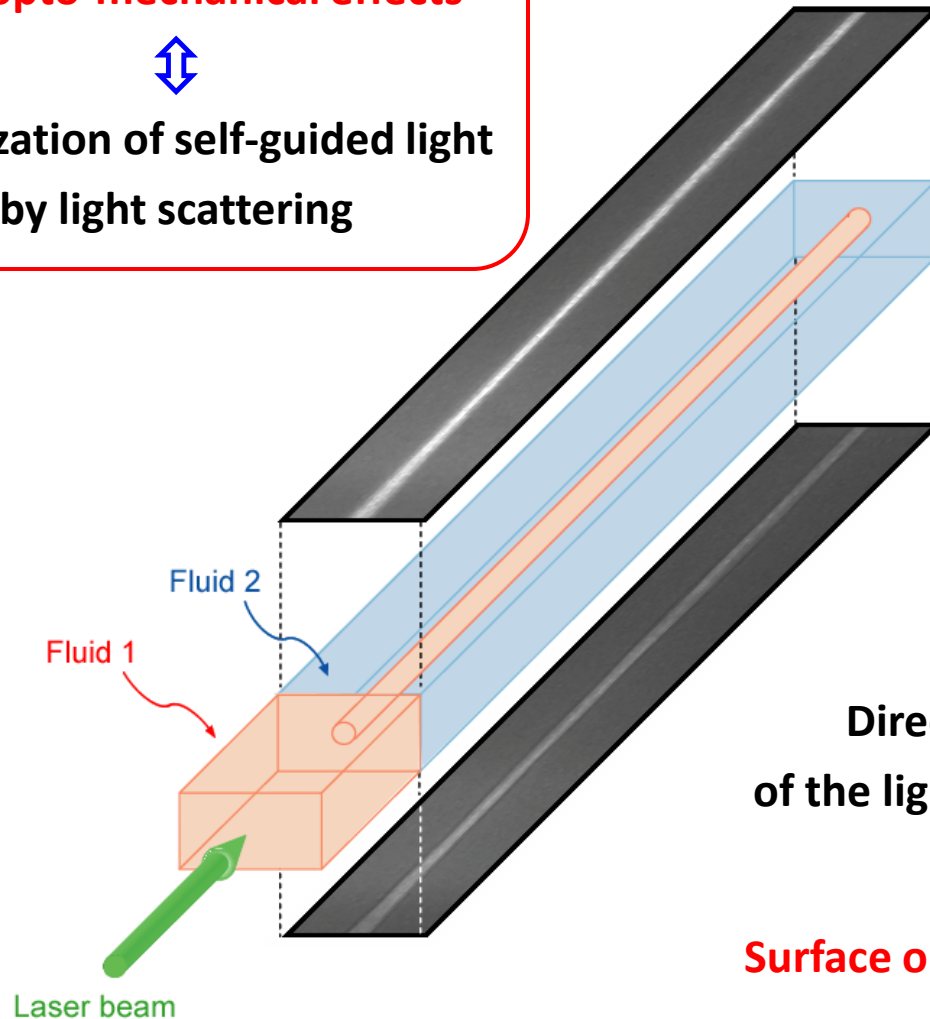


The micro-optical pipeline : how it looks in the lab

Bulk opto-mechanical effects



**Visualization of self-guided light
by light scattering**



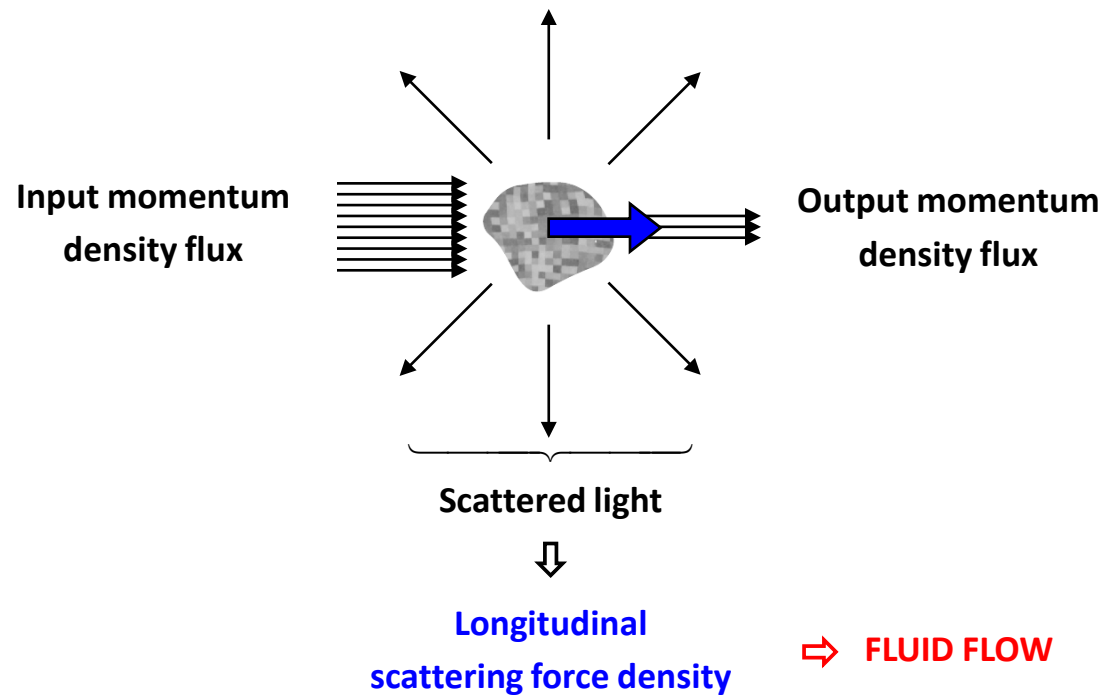
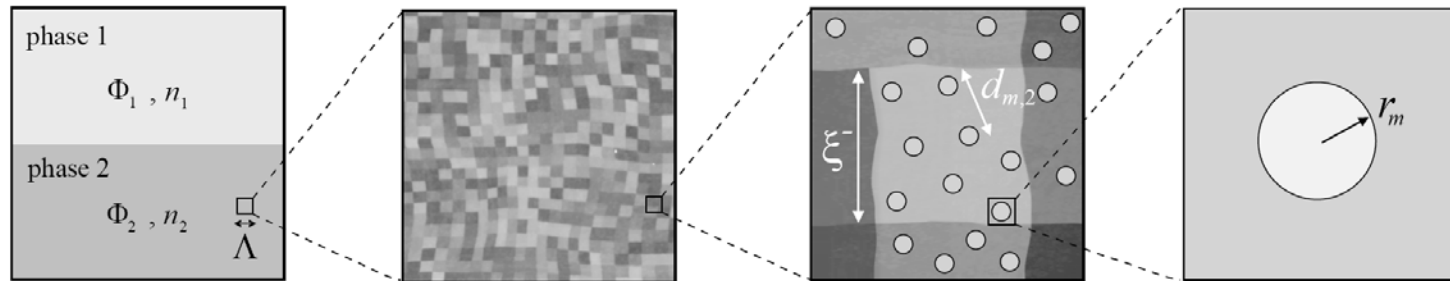
**Direct optical imaging
of the light-induced liquid fiber**



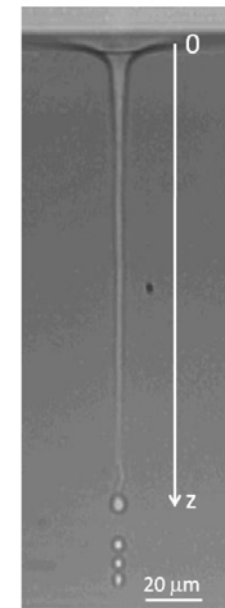
Surface opto-mechanical effects

Bulk contribution of the optical force density : light scattering

Micellar phases scatter light



Dripping jet

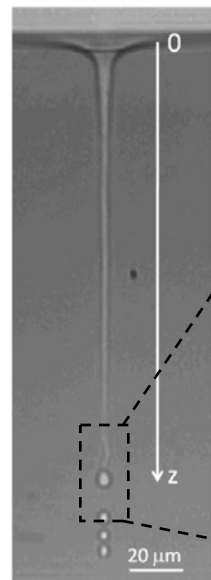


Light-induced bulk flow in the lab

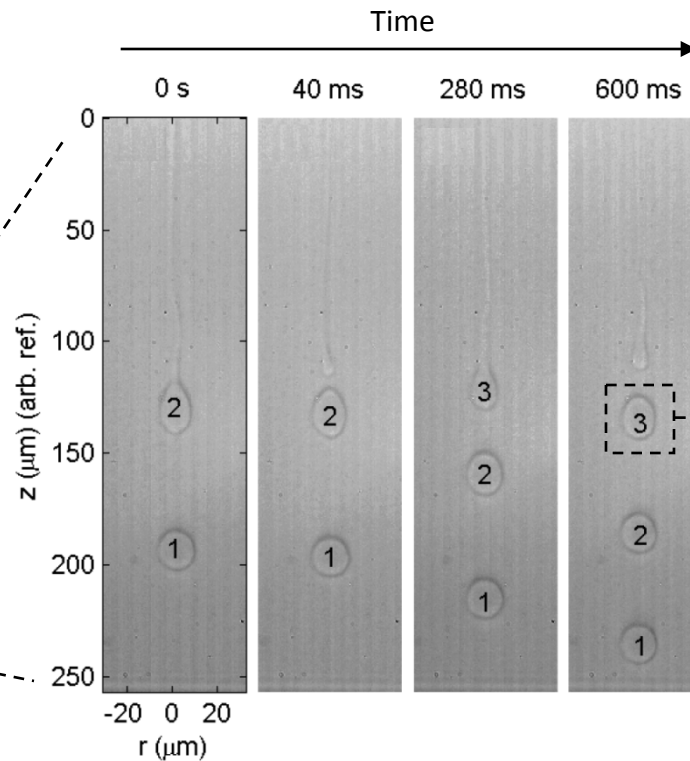
Scattering force



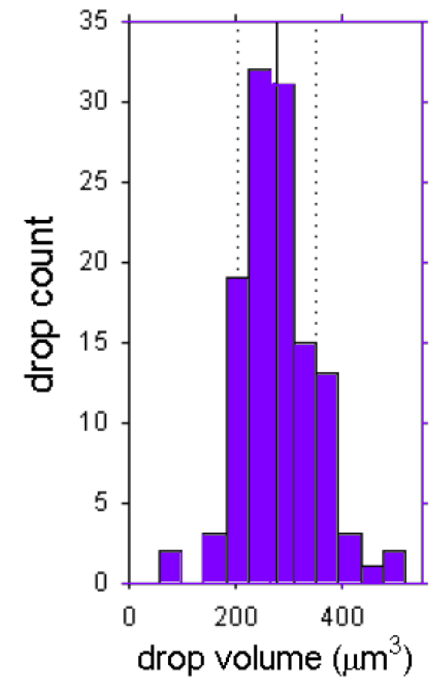
Dripping jet



$$w_0 = 3.5 \mu\text{m}$$

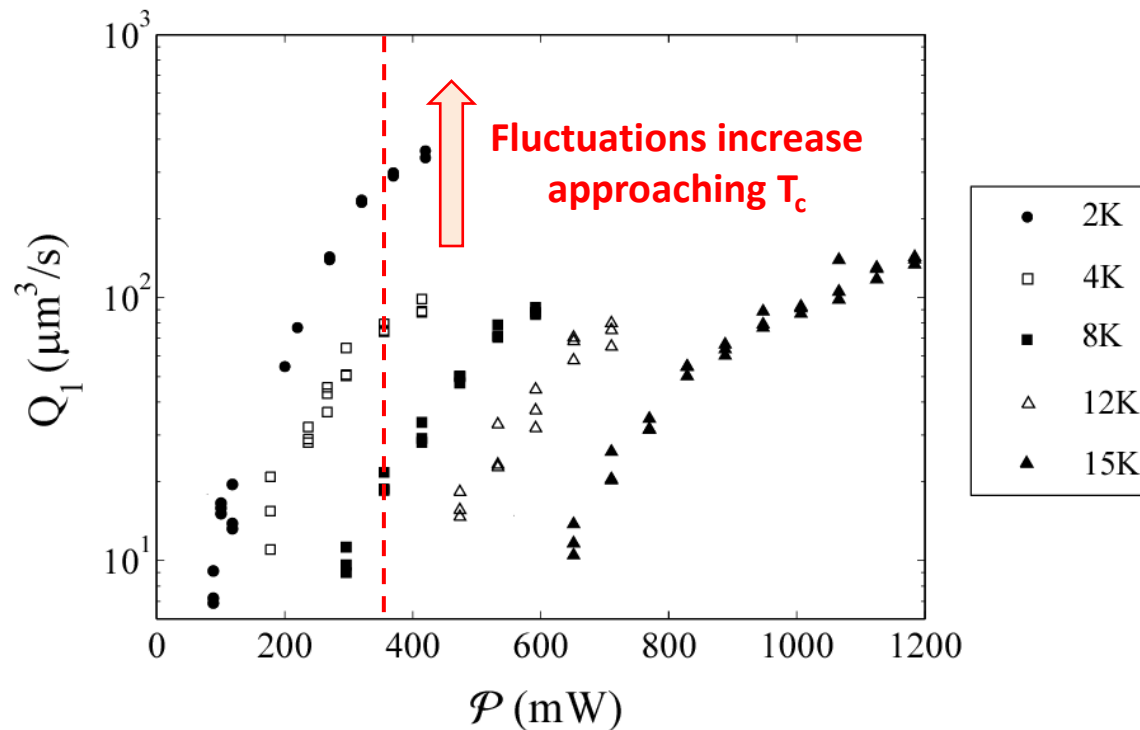


Flow rate statistics



Light-induced bulk flow in the lab

Average flow rate as a function of temperature and incident power



Light scattering is temperature dependent
via the size of refractive index fluctuations

Hydrodynamic model for light-induced flow

Bulk force balance

$$\mathbf{0} = \mathbf{f}_i^{\text{hydro}} + \mathbf{f}_i^{\text{em}} + \mathbf{f}_i^{\text{scatt}}, \quad i = 1, 2$$

$$\mathbf{f}_i^{\text{hydro}} = \nabla \cdot \mathbf{T}_i^{\text{hydro}}$$

$$\mathbf{T}_i^{\text{hydro}} = -p_i \mathbf{I} + \eta_i (\nabla \mathbf{v} + {}^t \nabla \mathbf{v})$$

$$\mathbf{f}_i^{\text{em}} = \nabla \cdot \mathbf{T}_i^{\text{em}}$$

$$\mathbf{T}_i^{\text{em}} = \frac{1}{2} \varepsilon_0 \rho_i \left. \frac{\partial \varepsilon_i}{\partial \rho_i} \right|_T \mathbf{E}^2 \mathbf{I} - \frac{1}{2} \varepsilon_0 \varepsilon_i \mathbf{E}^2 \mathbf{I} + \varepsilon_0 \varepsilon_i \mathbf{E} {}^t \mathbf{E}$$

Individual spherical scatterer

$$\mathbf{F}^{\text{scatt}} = \sigma n S / c$$

Micellar phases

$$\mathbf{f}_i^{\text{scatt}} = N \mathbf{F}_i^{\text{scatt}} = \tau_i n_i \mathbf{S}_i / c$$

$$\tau_i = \int \int \frac{1}{I_i} \frac{d^2 I_i}{d\Omega dz}(\Omega) d\Omega$$

Surface force balance

$$\left(T_2^{\text{hydro}} - T_1^{\text{hydro}} \right) \mathbf{n}_{1 \rightarrow 2} + \left(T_2^{\text{em}} - T_1^{\text{em}} \right) \mathbf{n}_{1 \rightarrow 2} - \sigma \kappa \mathbf{n}_{1 \rightarrow 2} = \mathbf{0}$$

Translational invariance along beam propagation

$$\mathbf{E}_i = \mathbf{E}_i(r)$$

$$\mathbf{f}_i^{\text{scatt}} = f_i^{\text{scatt}}(r)\hat{\mathbf{z}}$$

$$\mathbf{v}_i = v_i(r)\hat{\mathbf{z}}$$



Surface force balance

$$\eta_1 \frac{dv_{z,1}}{dr} \Big|_{r=R} = \eta_2 \frac{dv_{z,2}}{dr} \quad (\text{axial contribution})$$

$$\Pi_{\text{Laplace}} = \Pi_{\text{radiation}} \quad (\text{radial contribution})$$

Bulk force balance

$$\eta_i \Delta \mathbf{v}_i + \mathbf{f}_i^{\text{scatt}} = K$$

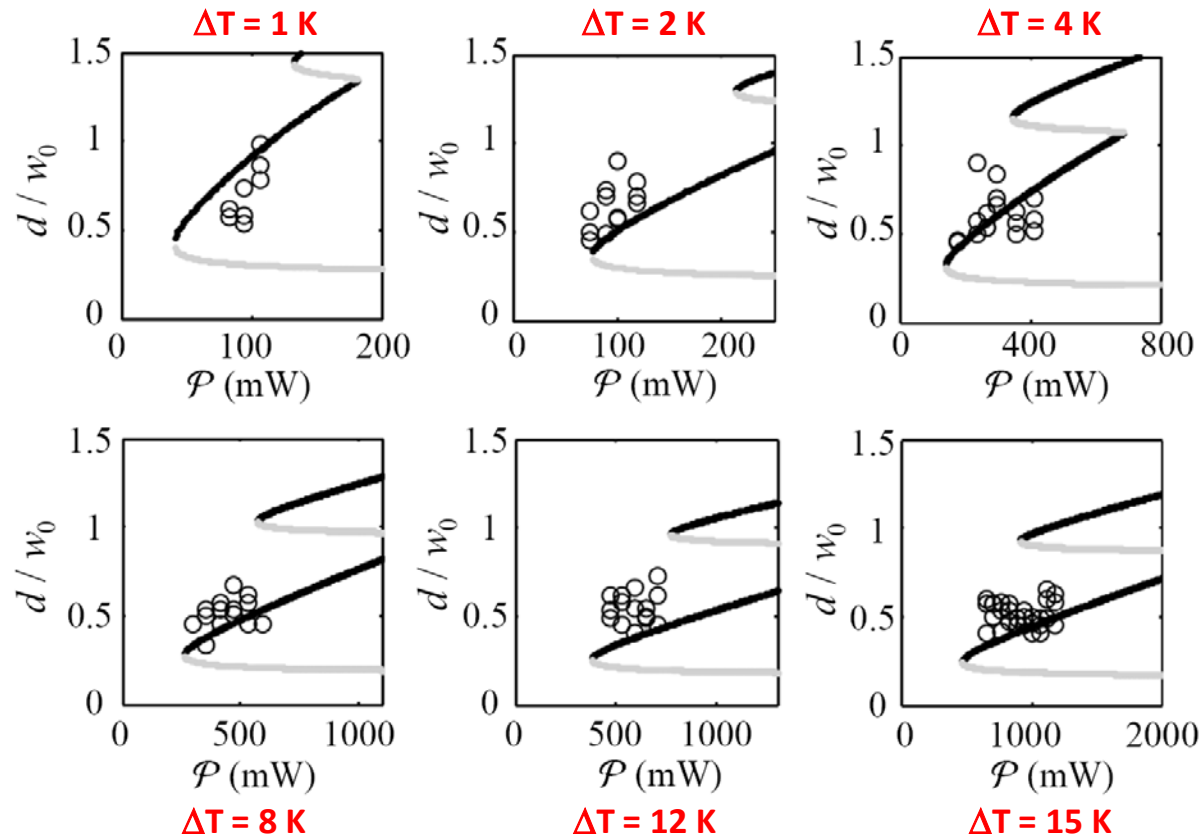
Fluid flow and light propagation inside micro-pipeline are decoupled

Geometry \Leftrightarrow waveguiding properties

Flow rate \Leftrightarrow scattering force density

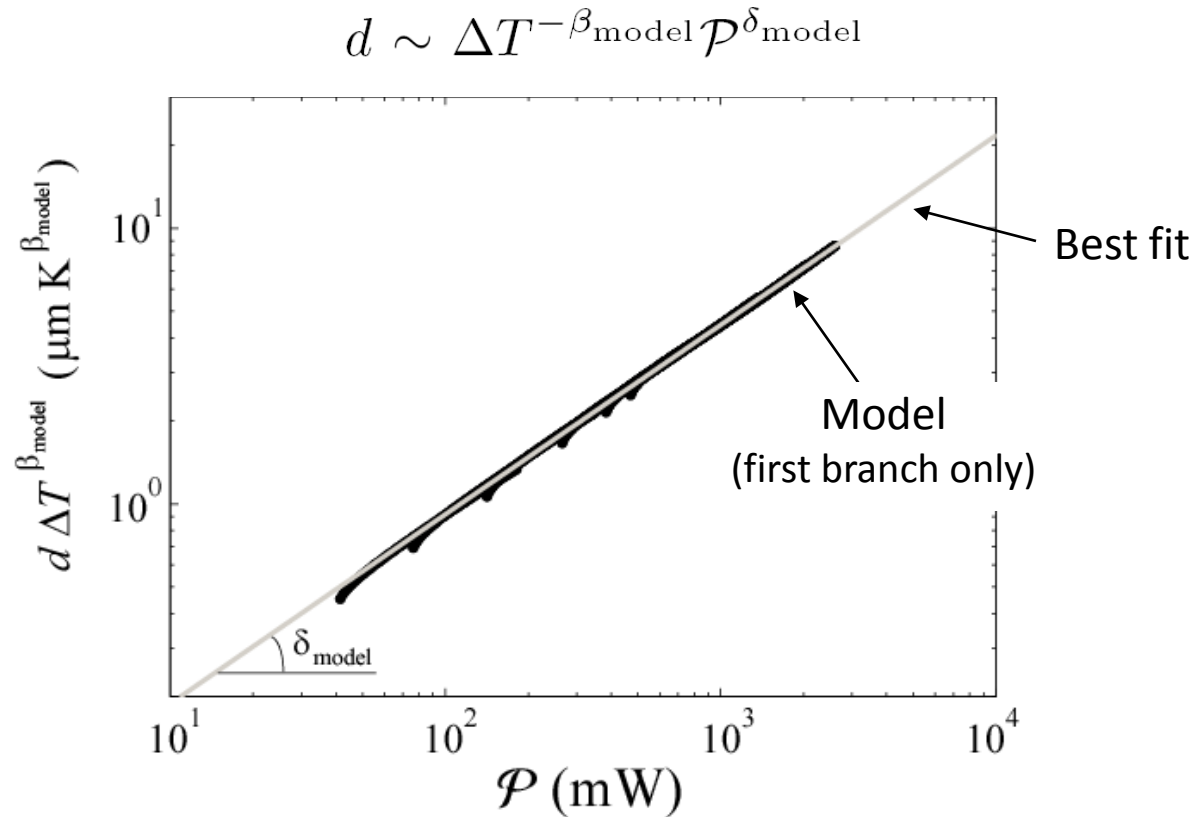
Optical micro-pipeline shape : experiment

$$\Delta T = T - T_c$$



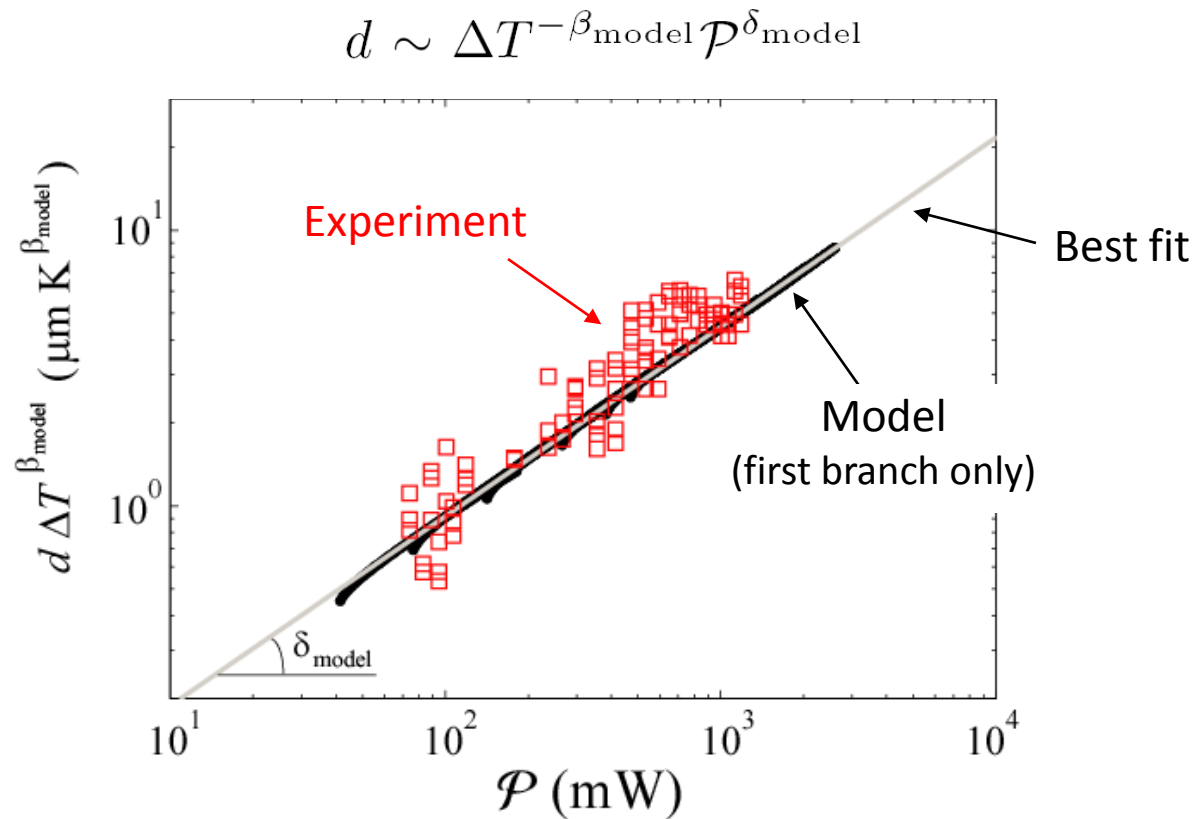
Only the first stable branch is explored for $w_0 = 3.5\mu\text{m}$

Optical micro-pipeline shape : model



Overall behavior is well-described by an empirical power law fit

Optical micro-pipeline shape : experiment / model



Satisfying agreement is obtained

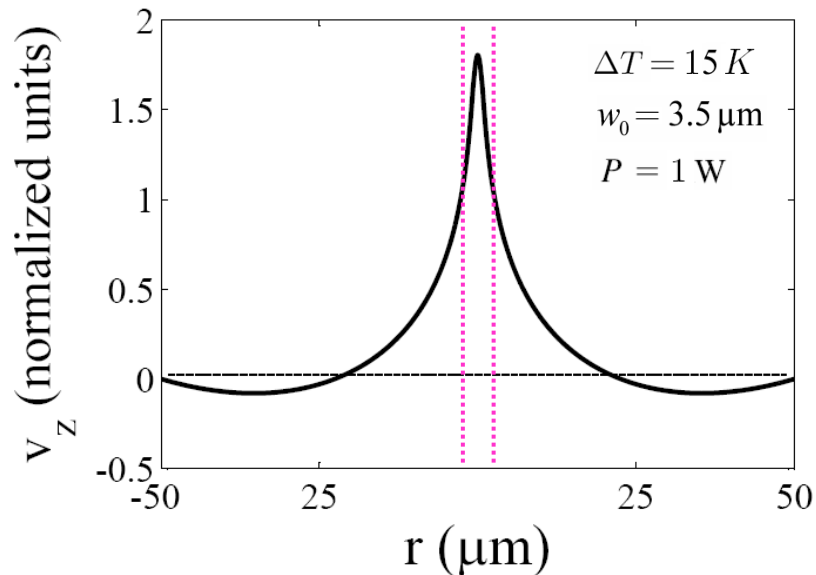
Optical micro-pipeline flow rate : model

Bulk force balance

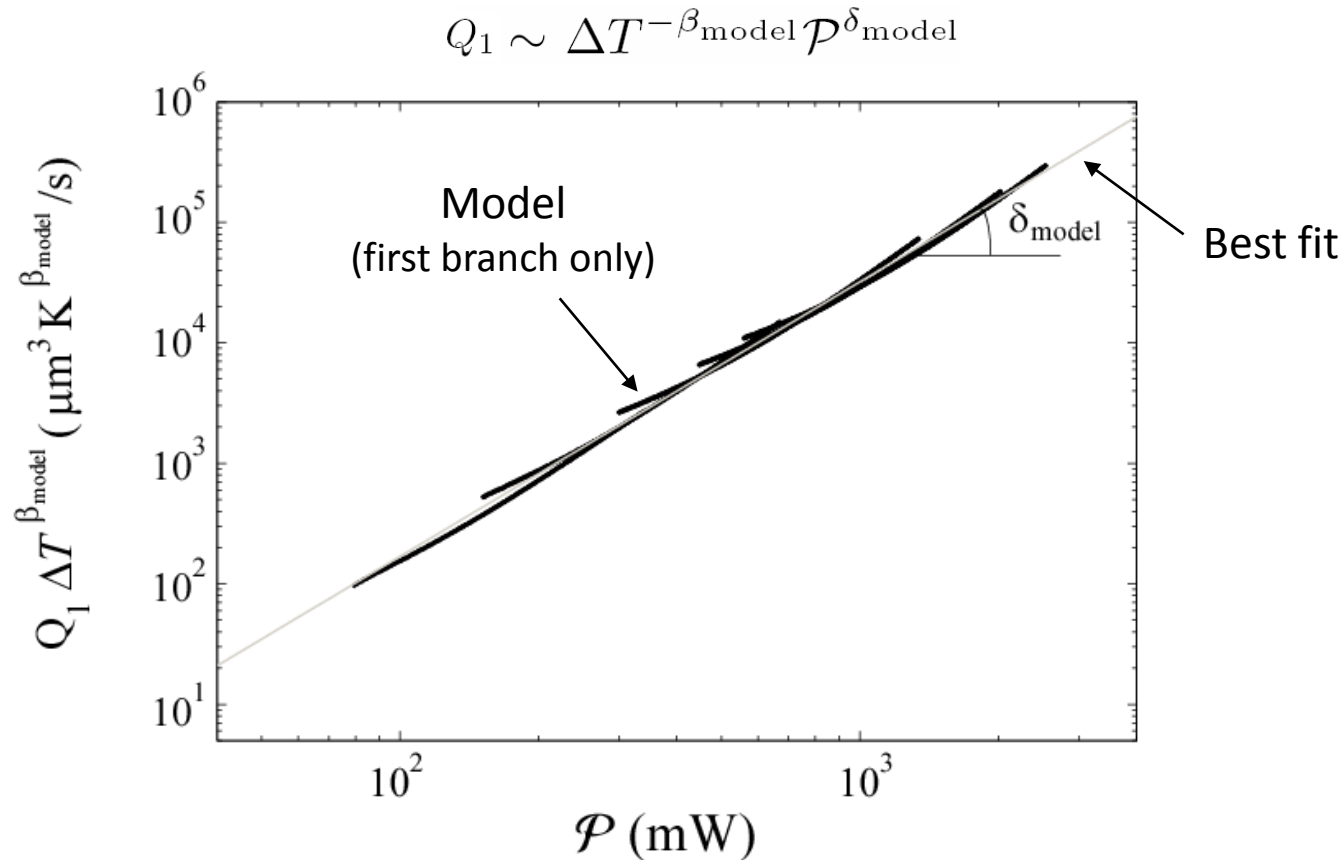
$$\eta_i \Delta \mathbf{v}_i + \mathbf{f}_i^{\text{scatt}} = K \quad \Rightarrow \quad \text{4 boundary conditions}$$

$$\left\{ \begin{array}{l} v_{z,1} = v_{z,2} \\ r=R \\ \eta_1 \frac{dv_{z,1}}{dr} = \eta_2 \frac{dv_{z,2}}{dr} \\ r=R \\ v_{z,2} = 0 \\ r=R_{\text{box}} \\ Q_2 = \int_R^{R_{\text{box}}} 2\pi r v_{z,2}(r) \, dr = 0 \end{array} \right.$$

Calculated velocity profile

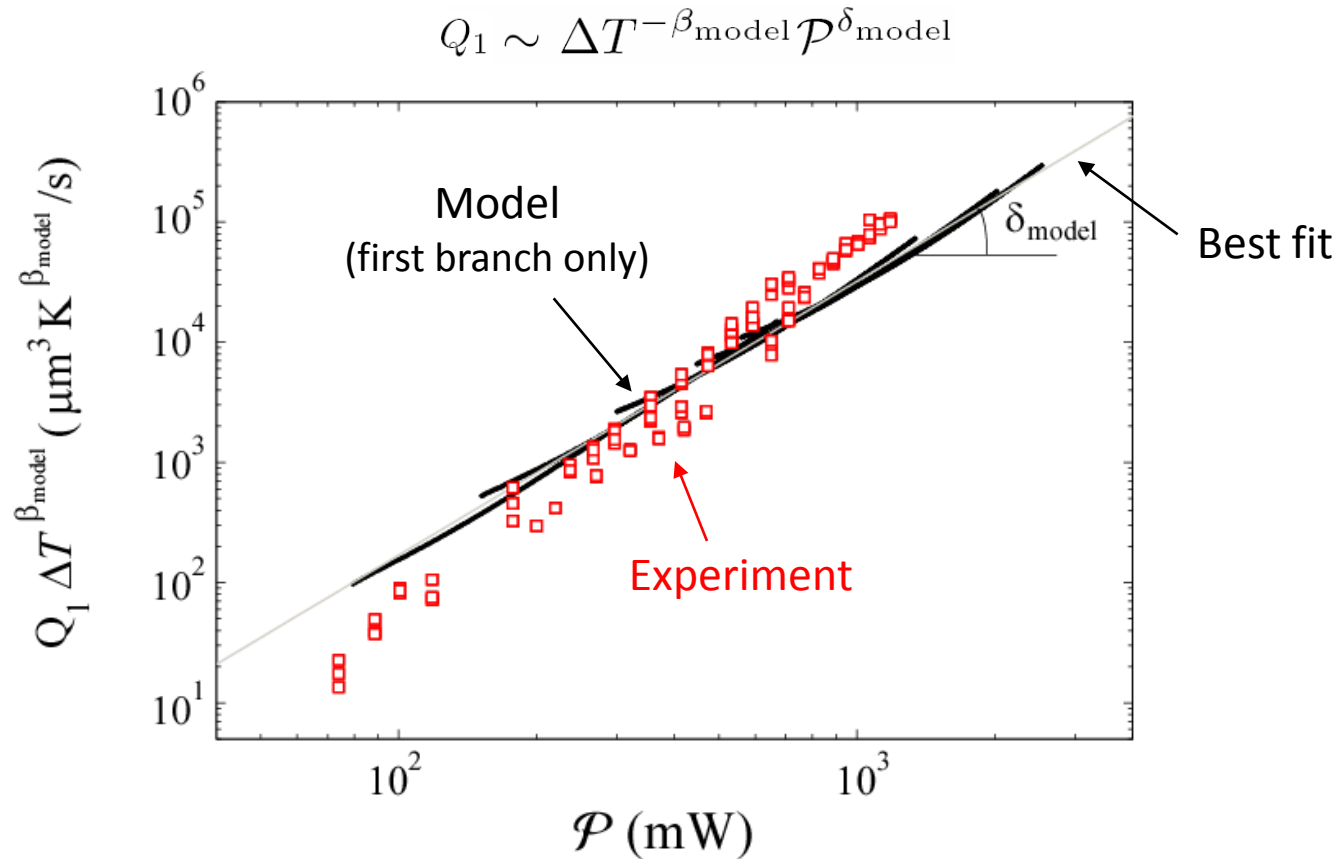


Optical micro-pipeline shape : experiment / model



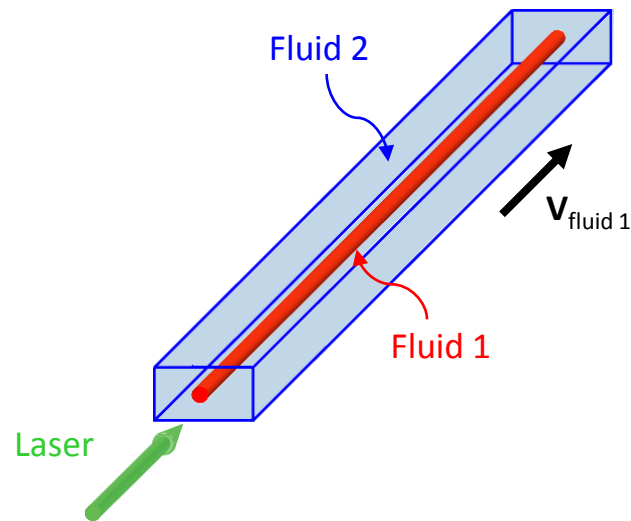
Overall behavior is well-described by an empirical power law fit

Optical micro-pipeline shape : experiment / model



Satisfying agreement is obtained

All-optical micro-pipeline / Fluid-mediated optical interconnect



Geometry

Light radiation pressure from propagating waveguided modes

Fluid flow

Light scattering force density

Experimental realization and quantitative theoretical description

E. Brasselet *et al.*, PRL **101**, 014501 (2008)

E. Brasselet *et al.*, PRA **78**, 013835 (2008)

R. Wunenburger *et al.*, J. Fluid Mech. **666**, 273 (2011)

Outline

1. Laser induced giant fluid interface deformations
2. Liquid-core liquid-cladding self-induced optical fibers
3. The optical micro-pipeline
4. **Acoustic analogies**

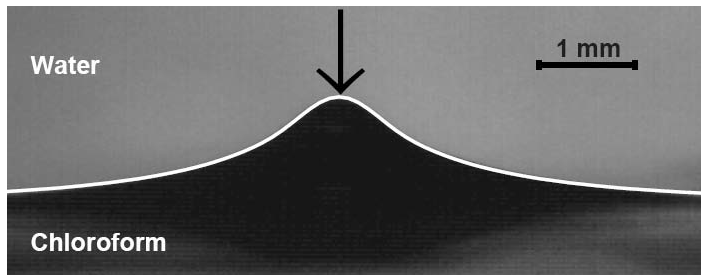
Mechanical effect of sound

Mechanical effects of waves are not restricted to light !

Acoustic radiation pressure

$$\Pi \sim I / c$$

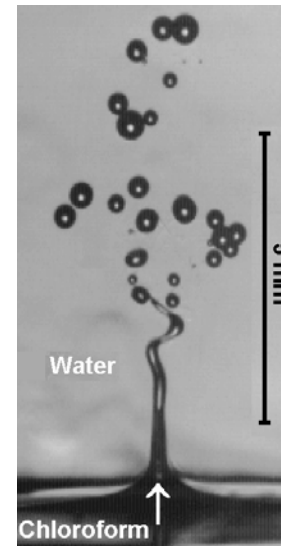
Acoustic fields are more efficient
than light fields ($\sim 10^5$)



B. Issenmann *et al.*, *Europhys. Lett.* **83**, 34002 (2008)

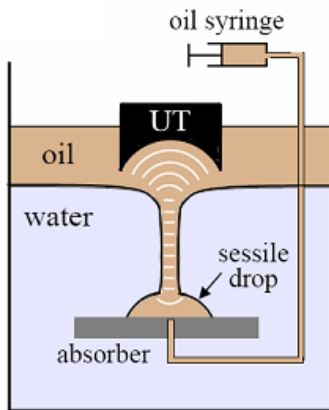
Acoustic radiation force density

Acoustic-matter interaction is dissipative
while it is conservative for light (in our case)

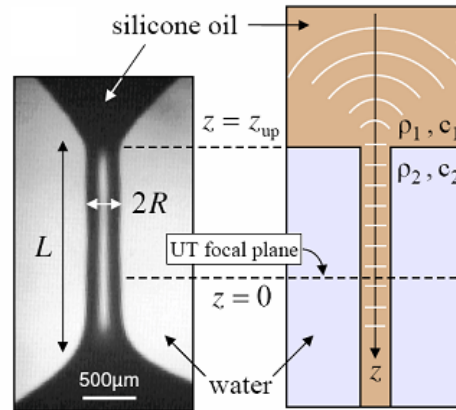


The acoustic milli-pipeline

The experiment

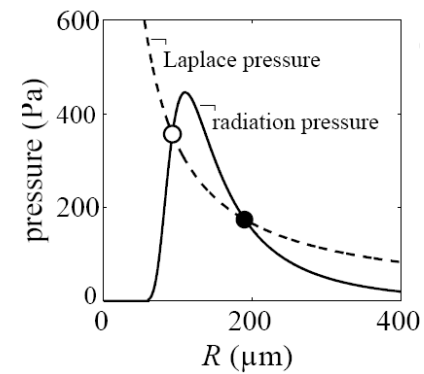


The model

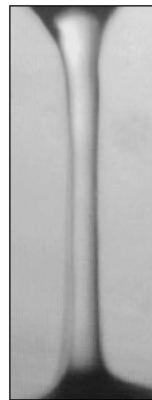


Acoustic waveguiding

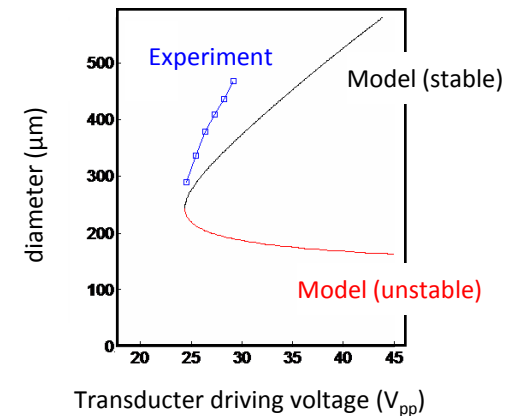
$$c_1 = 998 \text{ m.s}^{-1} < c_2 = 1488 \text{ m.s}^{-1}$$



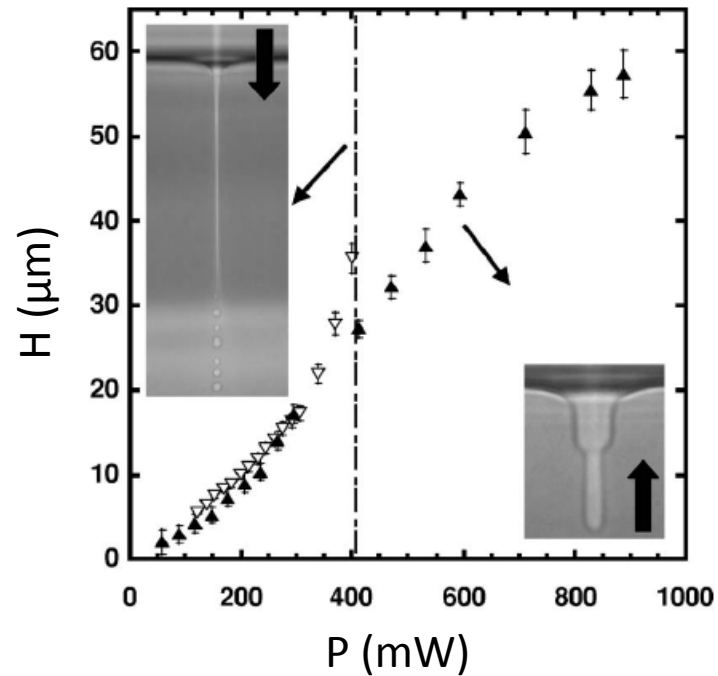
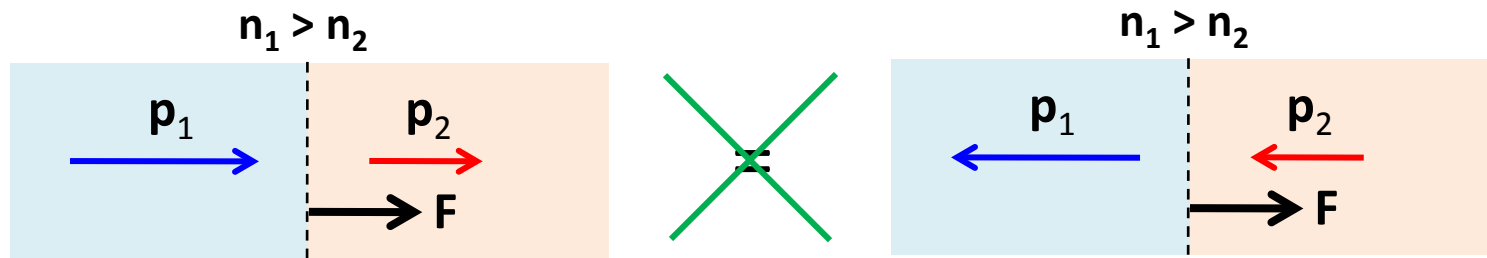
$$\Lambda \sim 10$$



N. Bertin *et al.*, PRL **105**, 164501 (2010)

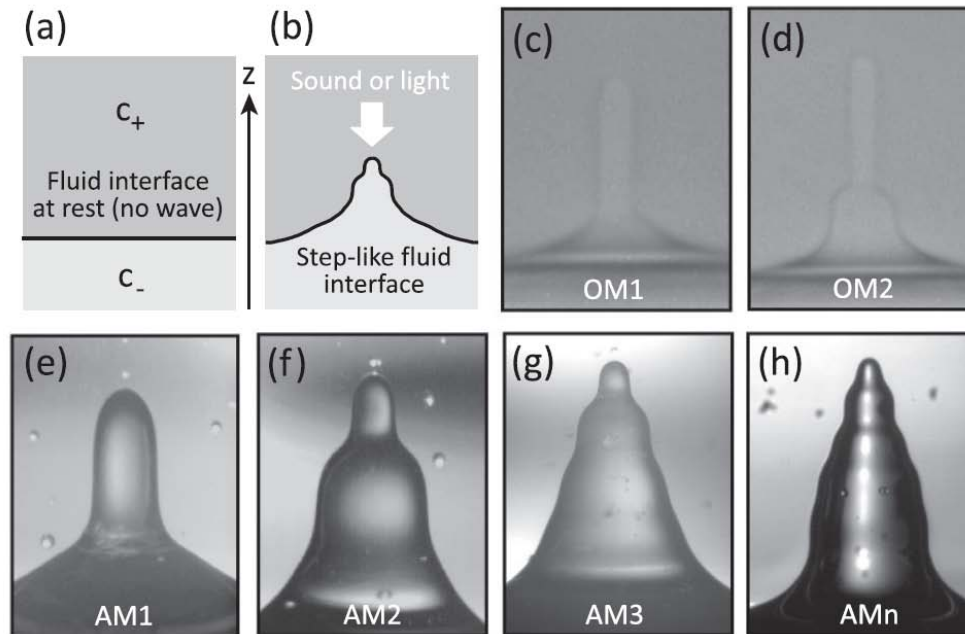


A closer look to up/ down symmetry



R. Wunenburger *et al.*, PRE **73**, 036314 (2006)

Analogy between optics and acoustics remains

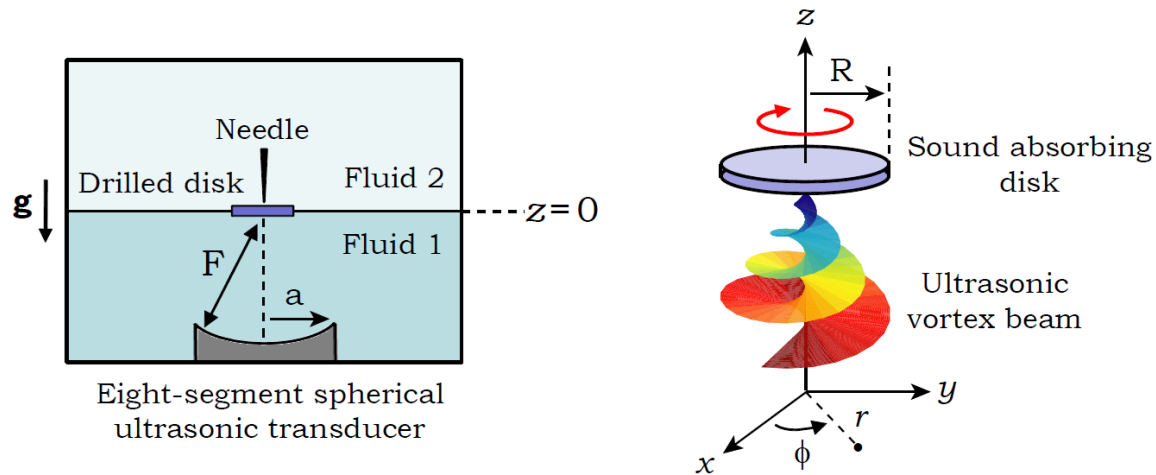


Universal morphologies of fluid interfaces deformed by the radiation pressure of acoustic or electromagnetic waves

N. Bertin *et al.*, PRL **109**, 244304 (2012)

Sound can also carry orbital angular momentum !

Rotational acousto-mechanics



A. Anhauser *et al.*, PRL **109**, 034301 (2012)